



Aplicação de Volumes Finitos a Modelo de Vaso de Adsorção

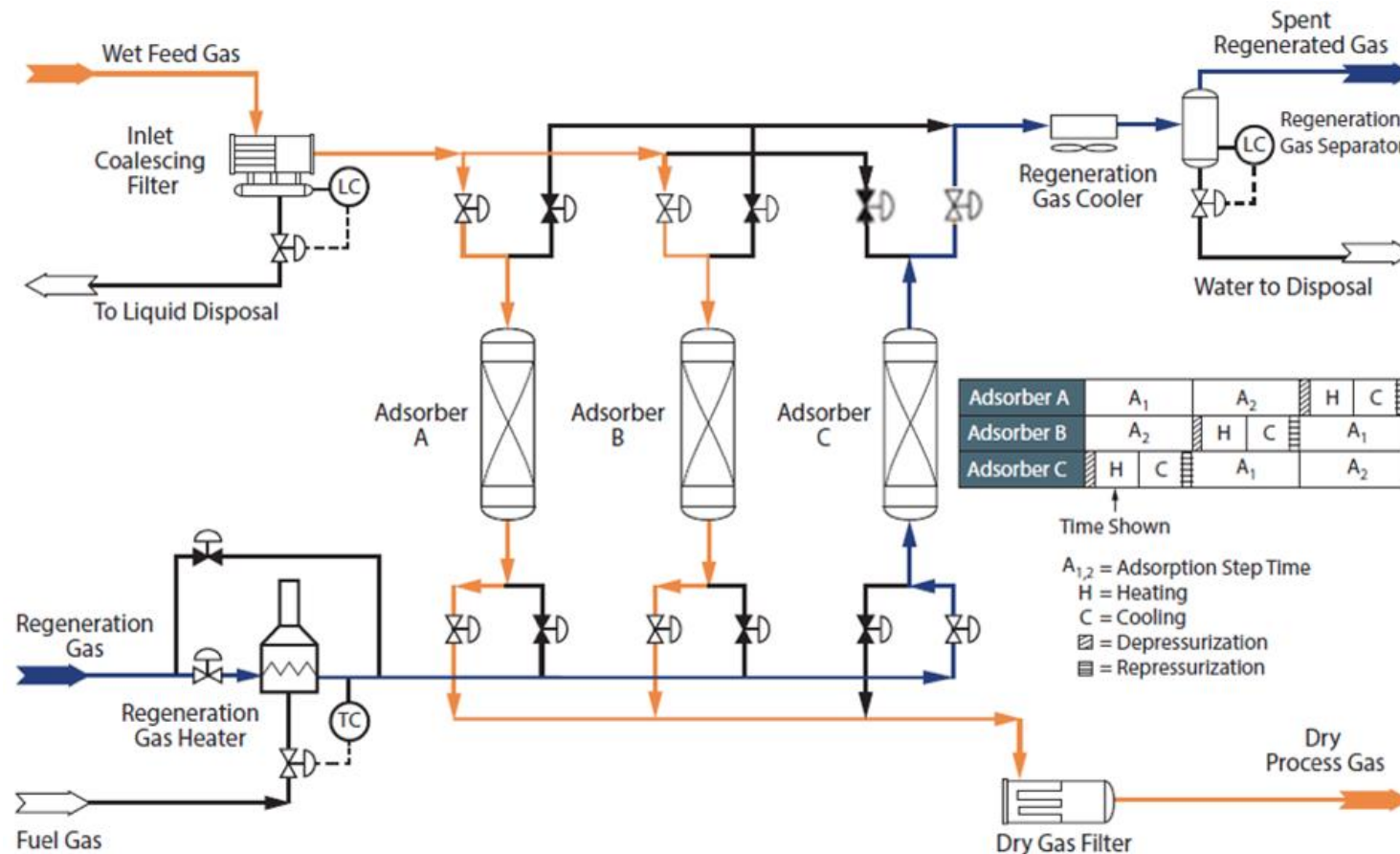
COQ-862 - Métodos Numéricos para Sistemas Distribuídos

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Desidratação de Gás Natural (*PPP offshore*)



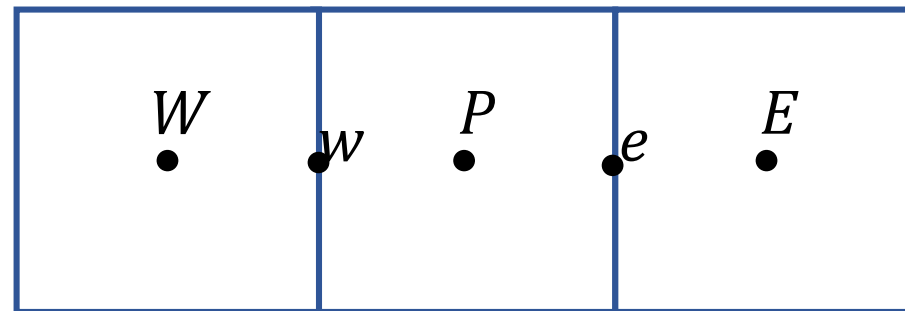
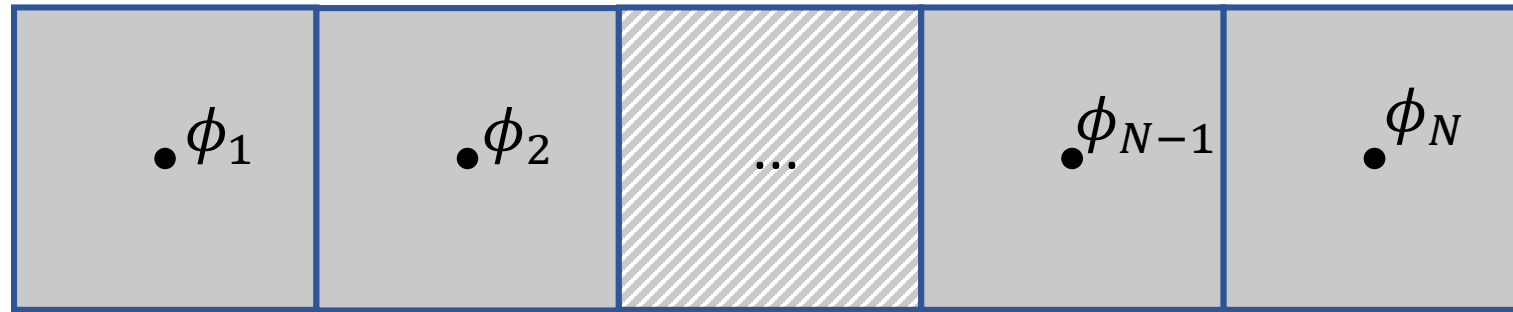
Campbell, J.M., *Gas Conditioning and Processing, Volume 2: The Equipment Modules, 9th Edition, 2nd Printing*, Editors Hubbard, R. and Snow-McGregor, K., Campbell Petroleum Series, Norman, Oklahoma, 2014.

Hipótesis Simplificadoras

<i>H1</i>	Somente dois componentes no gás: água e etanol. Considera-se que o etanol não é adsorvido pelas peneiras moleculares.
<i>H2</i>	Gás ideal
<i>H3</i>	Velocidade constante
<i>H4</i>	Sistema adiabático
<i>H5</i>	A resistência à transferência de massa segue um modelo linear (LDF)
<i>H6</i>	Isoterma de Langmuir
⋮	⋮

Volumes Finitos

Volumes Discretos de Mesmo Tamanho



Conservação de Massa da Fase Fluida

Modelo

$\frac{\partial c_a}{\partial t} = \frac{\partial^2(D_{ax} \cdot c_a)}{\partial z^2} - \frac{\partial\left(\frac{u}{\epsilon} \cdot c_a\right)}{\partial z} - \left(\frac{1-\epsilon}{\epsilon}\right) \rho_s \cdot \frac{\partial q_a}{\partial t}$	Concentração molar de água
$D_{ax} \cdot \frac{\partial c_a}{\partial z} \Big _{w=0} - \left(\frac{u}{\epsilon} \cdot c_a\right) \Big _{w=0} = -u_f \cdot c_{af}$	Condição de contorno em $z = 0$
$\frac{\partial c_a}{\partial z} \Big _{e=L} = 0$	Condição de contorno em $z = L$

Concentração em um volume interno P

$\int_w^e \frac{\partial c_a}{\partial t} dz = \int_w^e \frac{\partial^2 (D_{ax} \cdot c_a)}{\partial z^2} dz - \int_w^e \frac{\partial \left(\frac{u}{\epsilon} \cdot c_a \right)}{\partial z} dz - \int_w^e S_c(Y, T, P) dz$	
$\frac{\partial c_{aP}}{\partial t} \cdot \Delta z = D_{ax} \cdot \left(\frac{\partial c_a}{\partial z} \Big _e - \frac{\partial c_a}{\partial z} \Big _w \right) - \frac{1}{\epsilon} \cdot \left((u \cdot c_a) \Big _e - (u \cdot c_a) \Big _w \right) - \bar{S}_c \cdot \Delta z$	
$\frac{\partial c_{aP}}{\partial t} = \frac{D_{ax}}{\Delta z} \cdot \left(\frac{\partial c_a}{\partial z} \Big _e - \frac{\partial c_a}{\partial z} \Big _w \right) - \frac{1}{\epsilon \Delta z} \cdot \left((u \cdot c_a) \Big _e - (u \cdot c_a) \Big _w \right) - \bar{S}_c$	
$\frac{\partial c_{aP}}{\partial t} = \frac{D_{ax}}{\Delta z} \cdot \left(\frac{\partial c_a}{\partial z} \Big _e - \frac{\partial c_a}{\partial z} \Big _w \right) - \frac{1}{\epsilon \Delta z} \cdot \left((u \cdot c_a) \Big _e - (u \cdot c_a) \Big _w \right) - \bar{S}_c$	
$\frac{\partial c_{aP}}{\partial t} = \frac{D_{ax}}{\Delta z^2} \cdot (c_{aP+1} - 2c_{aP} + c_{aP-1}) - \frac{1}{\epsilon \Delta z} \cdot (u_P c_{aP} - u_{P-1} c_{aP-1}) - \bar{S}_c$	<p>Aproximação de primeira ordem para os termos advectivo e dispersivo.</p>

Concentração no primeiro volume

$\frac{\partial c_{a1}}{\partial t} = \frac{D_{ax}}{\Delta z} \cdot \left(\frac{\partial c_a}{\partial z} \Big _e - \frac{\partial c_a}{\partial z} \Big _0 \right) - \frac{1}{\epsilon \Delta z} \cdot \left((u \cdot c_a) \Big _e - (u \cdot c_a) \Big _0 \right) - \bar{S}_c$	$P = 1$
$\frac{\partial c_{a1}}{\partial t} = \frac{D_{ax}}{\Delta z} \cdot \frac{\partial c_a}{\partial z} \Big _e - \frac{D_{ax}}{\Delta z} \cdot \frac{\partial c_a}{\partial z} \Big _0 - \frac{1}{\epsilon \Delta z} \cdot (u \cdot c_a) \Big _e + \frac{1}{\epsilon \Delta z} \cdot (u \cdot c_a) \Big _0 - \bar{S}_c$	
$\frac{\partial c_{a1}}{\partial t} = \frac{D_{ax}}{\Delta z} \cdot \frac{\partial c_a}{\partial z} \Big _e - \left(\frac{D_{ax}}{\Delta z} \cdot \frac{\partial c_a}{\partial z} \Big _0 - \frac{1}{\epsilon \Delta z} \cdot (u \cdot c_a) \Big _0 \right) - \frac{1}{\epsilon \Delta z} \cdot (u \cdot c_a) \Big _e - \bar{S}_c$	
$D_{ax} \cdot \frac{\partial c_a}{\partial z} \Big _0 - \frac{u}{\epsilon} \cdot c_a \Big _0 = - \left(\frac{F_f \cdot W_a}{MM_a} \right) = -u_f \cdot c_{af}$	Condição de contorno em $z = 0$
$\frac{\partial c_{a1}}{\partial t} = \frac{D_{ax}}{\Delta z} \cdot \frac{\partial c_a}{\partial z} \Big _e + (u_f \cdot c_{af}) - \frac{1}{\epsilon \Delta z} \cdot (u \cdot c_a) \Big _1 - \bar{S}_c$	Aproximação de primeira ordem para o termo advectivo
$\frac{\partial c_{a1}}{\partial t} = \frac{D_{ax}}{\Delta z^2} \cdot (c_{a2} - c_{a1}) + (u_f \cdot c_{af}) - \frac{1}{\epsilon \Delta z} \cdot (u \cdot c_a) \Big _1 - \bar{S}_c$	Aproximação de primeira ordem para o termo dispersivo

Concentração no último volume

$\frac{\partial c_{a_N}}{\partial t} = \frac{D_{ax}}{\Delta z} \cdot \left(\frac{\partial c_a}{\partial z} \Big _e - \frac{\partial c_a}{\partial z} \Big _w \right) - \frac{1}{\epsilon \Delta z} \cdot \left[(u \cdot c_a) \Big _e - (u \cdot c_a) \Big _w \right] - \bar{S}_c$	$P = N$
$\frac{\partial c_a}{\partial z} \Big _e = 0$	Condição de contorno em $z = L$
$\frac{\partial c_{a_N}}{\partial t} = \frac{D_{ax}}{\Delta z} \cdot \left(0 - \frac{\partial c_a}{\partial z} \Big _w \right) - \frac{1}{\epsilon \Delta z} \cdot \left[(u \cdot c_a) \Big _e - (u \cdot c_a) \Big _w \right] - \bar{S}_c$	
$\frac{\partial c_{a_N}}{\partial t} = -\frac{D_{ax}}{\Delta z} \cdot \frac{\partial c_a}{\partial z} \Big _w - \frac{1}{\epsilon \Delta z} \cdot \left[(u \cdot c_a) \Big _e - (u \cdot c_a) \Big _w \right] - \bar{S}_c$	
$\frac{\partial c_{a_N}}{\partial t} = -\frac{D_{ax}}{\Delta z} \cdot \frac{\partial c_a}{\partial z} \Big _w - \frac{1}{\epsilon \Delta z} \cdot \left[(u \cdot c_a) \Big _N - (u \cdot c_a) \Big _{N-1} \right] - \bar{S}_c$	Aproximação de primeira ordem pra trás nos termos advectivos
$\frac{\partial c_{a_N}}{\partial t} = -\frac{D_{ax}}{\Delta z^2} \cdot (c_{a_N} - c_{a_{N-1}}) - \frac{1}{\epsilon \Delta z} \cdot \left[(u \cdot c_a) \Big _N - (u \cdot c_a) \Big _{N-1} \right] - \bar{S}_c$	Aproximação de primeira ordem centrais nos termos dispersivos

Conservação de Massa no Sólido

$$\frac{\partial q}{\partial t} = K_{LDF} \cdot [q_{eq}(z) - q]$$

$$\int_w^e \frac{\partial q}{\partial t} dz = K_{LDF} \int_w^e q_{eq}(z) dz - K_{LDF} \int_w^e q dz$$

$$\frac{\partial q_P}{\partial t} \Delta Z = K_{LDF} \cdot q_{eq_P} \cdot \Delta Z - K_{LDF} \cdot q_P \cdot \Delta Z$$

$$\frac{\partial q_P}{\partial t} = K_{LDF} \cdot (q_{eq_P} - q_P)$$

Conservação da Massa Global

$\frac{\partial C}{\partial t} = -\frac{1}{\epsilon} \cdot \frac{\partial(u \cdot C)}{\partial z} - \left(\frac{1-e}{e}\right) \cdot \rho_s \cdot \frac{\partial q}{\partial t}$	
$(u \cdot C) \Big _{w=0} = \frac{F_f}{A \cdot MM_{tot}}$	Condição de contorno em $z = 0$.
$\frac{\partial C}{\partial t} = -\frac{1}{\epsilon \Delta z} \left[(u \cdot C) \Big _e - (u \cdot C) \Big _w \right] - \bar{S}_C(Y, T, P)$	Equação em termos de fluxo
$\frac{\partial C_P}{\partial t} = -\frac{1}{\epsilon} \left[(u \cdot C) \Big _P - (u \cdot C) \Big _{P-1} \right] - \bar{S}_{C_P}(Y, T, P)$	Aproximação de primeira ordem para os termos advectivos
$\frac{\partial C_1}{\partial t} = -\frac{1}{\epsilon} \left[(u \cdot C) \Big _1 - (u \cdot C) \Big _0 \right] - \bar{S}_{C_P}(Y, T, P)$	
$\frac{\partial C_1}{\partial t} = -\frac{1}{\epsilon} \left[(u \cdot C) \Big _1 - \left(\frac{F_f}{A \cdot MM_{tot}} \right) \right] - \bar{S}_{C_P}(Y, T, P)$	

Conservação de Energia

Conservação de Energia

$$\left[\rho_g c_{\rho_g} + \left(\frac{1 - \epsilon}{\epsilon} \right) \rho_s c_{\rho_s} \right] \frac{\partial T}{\partial t} = \frac{\partial^2 (k_{ef} \cdot T)}{\partial z^2} - \frac{\partial \left(\frac{u \rho_g c_{\rho_g}}{\epsilon} \cdot T \right)}{\partial z} - Q \cdot \left(\frac{1 - \epsilon}{\epsilon} \right) \rho_s \cdot \frac{\partial q}{\partial t}$$

$$\frac{k_{ef}}{\left[\rho_g c_{\rho_g} + \left(\frac{1 - \epsilon}{\epsilon} \right) \rho_s c_{\rho_s} \right]} = \alpha$$

$$\frac{u \rho_g c_{\rho_g}}{\epsilon \left[\rho_g c_{\rho_g} + \left(\frac{1 - \epsilon}{\epsilon} \right) \rho_s c_{\rho_s} \right]} = \beta$$

$$\frac{Q \cdot \left(\frac{1 - \epsilon}{\epsilon} \right) \rho_s}{\left[\rho_g c_{\rho_g} + \left(\frac{1 - \epsilon}{\epsilon} \right) \rho_s c_{\rho_s} \right]} = \gamma$$

$$\frac{\partial T}{\partial t} = \frac{\partial^2 (\alpha \cdot T)}{\partial z^2} - \frac{\partial (\beta \cdot T)}{\partial z} - \gamma \cdot \frac{\partial q}{\partial t}$$

Temperatura em um volume interno ($P = 2:N-1$)

$\frac{\partial T}{\partial t} = \frac{\partial^2(\alpha \cdot T)}{\partial z^2} - \frac{\partial(\beta \cdot T)}{\partial z} - \gamma \cdot \frac{\partial q}{\partial t}$	Assumindo que a densidade varia pouco com o tempo e com z.
$\frac{\partial T_P}{\partial t} \Delta z = \alpha \left(\frac{\partial T}{\partial z} \Big _e - \frac{\partial T}{\partial z} \Big _w \right) - \beta \left(T \Big _e - T \Big _w \right) - \gamma \cdot \frac{\partial q_P}{\partial t} \Delta z$	
$\frac{\partial T_P}{\partial t} = \frac{\alpha}{\Delta z} \left(\frac{\partial T}{\partial z} \Big _e - \frac{\partial T}{\partial z} \Big _w \right) - \frac{\beta}{\Delta z} \left(T \Big _e - T \Big _w \right) - \gamma \cdot \frac{\partial q_P}{\partial t}$	
$\frac{\partial T_P}{\partial t} = \frac{\alpha}{\Delta z^2} (T_{P+1} - 2T_P + T_{P-1}) - \frac{\beta}{\Delta z} (T_P - T_{P-1}) - \gamma \cdot \frac{\partial q_P}{\partial t}$	Aproximação de primeira ordem para os termos advectivo e dispersivo.

Temperatura no primeiro volume ($P = 1$)

$\alpha \cdot \frac{\partial T}{\partial z} \Big _0 - \beta \cdot T \Big _0 = -\beta \cdot \epsilon \cdot T_f$	Assumindo que a densidade varia pouco com o tempo e com z.
$\frac{\partial T_1}{\partial t} = \frac{\alpha}{\Delta z} \left(\frac{\partial T}{\partial z} \Big _e - \frac{\partial T}{\partial z} \Big _0 \right) - \frac{\beta}{\Delta z} \left(T \Big _e - T \Big _0 \right) - \gamma \cdot \frac{\partial q_1}{\partial t}$	$P = 1$
$\frac{\partial T_1}{\partial t} = \frac{1}{\Delta z} \cdot \left(\alpha \cdot \frac{\partial T}{\partial z} \Big _e - \beta \cdot T \Big _e \right) - \frac{1}{\Delta z} \left(\alpha \cdot \frac{\partial T}{\partial z} \Big _0 - \beta \cdot T \Big _0 \right) - \gamma \cdot \frac{\partial q_1}{\partial t}$	
$\frac{\partial T_1}{\partial t} = \frac{1}{\Delta z} \cdot \left(\alpha \cdot \frac{\partial T}{\partial z} \Big _e - \beta \cdot T \Big _e \right) + \frac{1}{\Delta z} (\beta \cdot \epsilon \cdot T_f) - \gamma \cdot \frac{\partial q_1}{\partial t}$	
$\frac{\partial T_1}{\partial t} = \frac{1}{\Delta z^2} \cdot [\alpha \cdot (T_{P+1} - T_P) - \beta \cdot T_P] + \frac{1}{\Delta z} (\beta \cdot \epsilon \cdot T_f) - \gamma \cdot \frac{\partial q_1}{\partial t}$	

Temperatura no primeiro volume (P = 1)

$\alpha \cdot \frac{\partial T}{\partial z} \Big _0 - \beta \cdot T \Big _0 = -\beta \cdot \epsilon \cdot T_f$	Assumindo que a densidade varia pouco com o tempo e com z.
$\frac{\partial T_1}{\partial t} = \frac{\alpha}{\Delta z} \left(\frac{\partial T}{\partial z} \Big _e - \frac{\partial T}{\partial z} \Big _0 \right) - \frac{\beta}{\Delta z} \left(T \Big _e - T \Big _0 \right) - \gamma \cdot \frac{\partial q_1}{\partial t}$	P = 1
$\frac{\partial T_1}{\partial t} = \frac{1}{\Delta z} \cdot \left(\alpha \cdot \frac{\partial T}{\partial z} \Big _e - \beta \cdot T \Big _e \right) - \frac{1}{\Delta z} \left(\alpha \cdot \frac{\partial T}{\partial z} \Big _0 - \beta \cdot T \Big _0 \right) - \gamma \cdot \frac{\partial q_1}{\partial t}$	
$\frac{\partial T_1}{\partial t} = \frac{1}{\Delta z} \cdot \left(\alpha \cdot \frac{\partial T}{\partial z} \Big _e - \beta \cdot T \Big _e \right) + \frac{1}{\Delta z} (\beta \cdot \epsilon \cdot T_f) - \gamma \cdot \frac{\partial q_1}{\partial t}$	
$\frac{\partial T_1}{\partial t} = \frac{1}{\Delta z^2} \cdot [\alpha \cdot (T_{P+1} - T_P) - \beta \cdot T_P] + \frac{1}{\Delta z} (\beta \cdot \epsilon \cdot T_f) - \gamma \cdot \frac{\partial q_1}{\partial t}$	

Temperatura no último volume (P = N)

$\frac{\partial T}{\partial z} \Big _e = 0$	P = N
$\frac{\partial T_N}{\partial t} = \frac{\alpha}{\Delta z} \left(\frac{\partial T}{\partial z} \Big _e - \frac{\partial T}{\partial z} \Big _w \right) - \frac{\beta}{\Delta z} (T \Big _N - T \Big _{N-1}) - \gamma \cdot \frac{\partial q_N}{\partial t}$	
$\frac{\partial T_N}{\partial t} = \frac{\alpha}{\Delta z} \left(0 - \frac{\partial T}{\partial z} \Big _w \right) - \frac{\beta}{\Delta z} (T \Big _N - T \Big _{N-1}) - \gamma \cdot \frac{\partial q_N}{\partial t}$	
$\frac{\partial T_N}{\partial t} = -\frac{\alpha}{\Delta z^2} (T_N - T_{N-1}) - \frac{\beta}{\Delta z} (T \Big _N - T \Big _{N-1}) - \gamma \cdot \frac{\partial q_N}{\partial t}$	

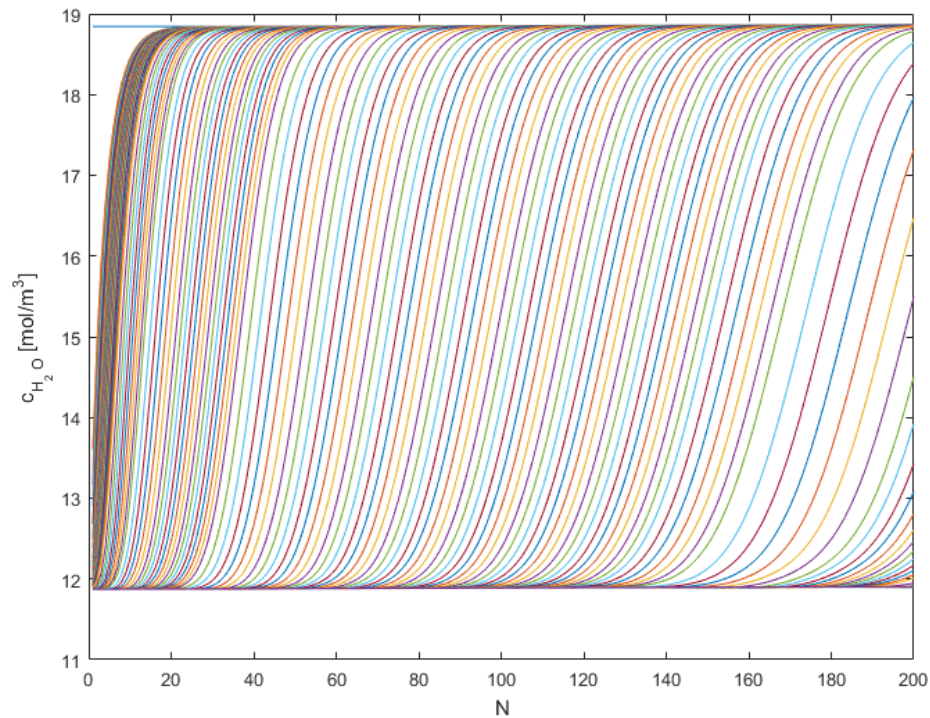
Simulação

Desidratação de Etanol

<i>Parâmetro e Constantes</i>	<i>Valor</i>	<i>Unidade</i>	<i>Parâmetro e Constantes</i>	<i>Valor</i>	<i>Unidade</i>
MM _{H₂O}	18	g	T _b	450	K
MM _{C₂H₆O}	46	g	N	300	pontos
R	8.314598	$J(K\ mol)^{-1}$	dL	0.0244	m
MF _w	0.182	%	ρ _s	729	kg/m ³
c _{ps}	1260	$J(kg\ K)^{-1}$	r _p	0.0015875	m
F _f	5.6694 X 100	kg/s	q _s	10.6656	mol/kg
T _f	440	K	T ₀	323	K
P _f	379.2120	kPa	K ₀	0.0441765	Pa^{-1}
L	7.3	m	K _{LDF}	0.5E-4	
D	2.4	m	μ	1e-6	
A	4.5239	m ²	RelTol	2e-4	
ε	0.63	[]	AbsTol	3e-3	
ε _b	0.4	[]	D _{ax}	0.00113	m^2s^{-1}
Q	51900	J	k _{ef}	41.26	$W(m\ K)^{-1}$
c _{pg}	1000	$J(kg\ K)^{-1}$			

M. Simo, C. J. Brown, and V. Hlavacek, "Simulation of pressure swing adsorption in fuel ethanol production process," *Comput. Chem. Eng.*, vol. 32, no. 7, pp. 1635–1649, 2008.

Concentração de Água no Gás

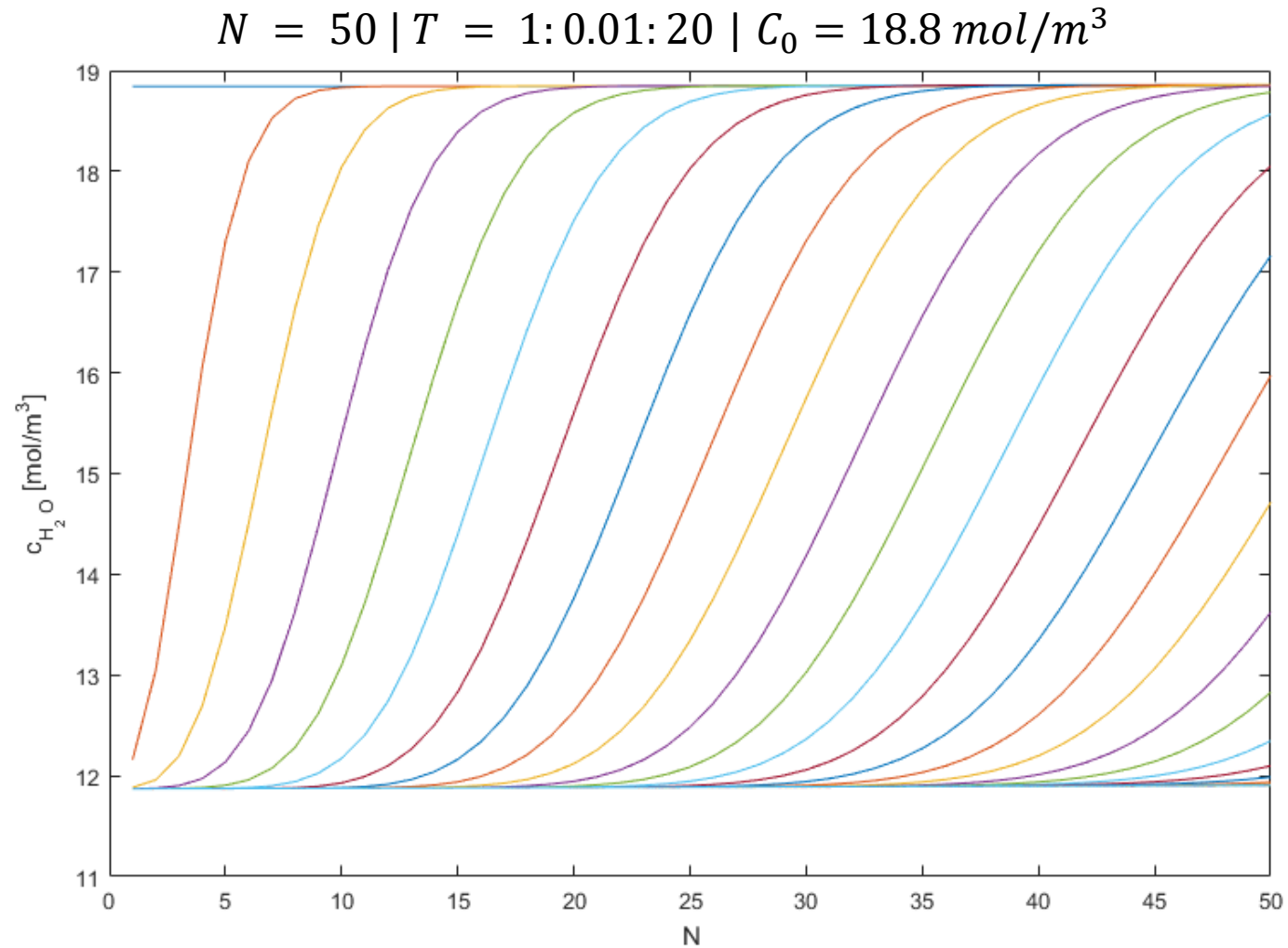


$N = 200$

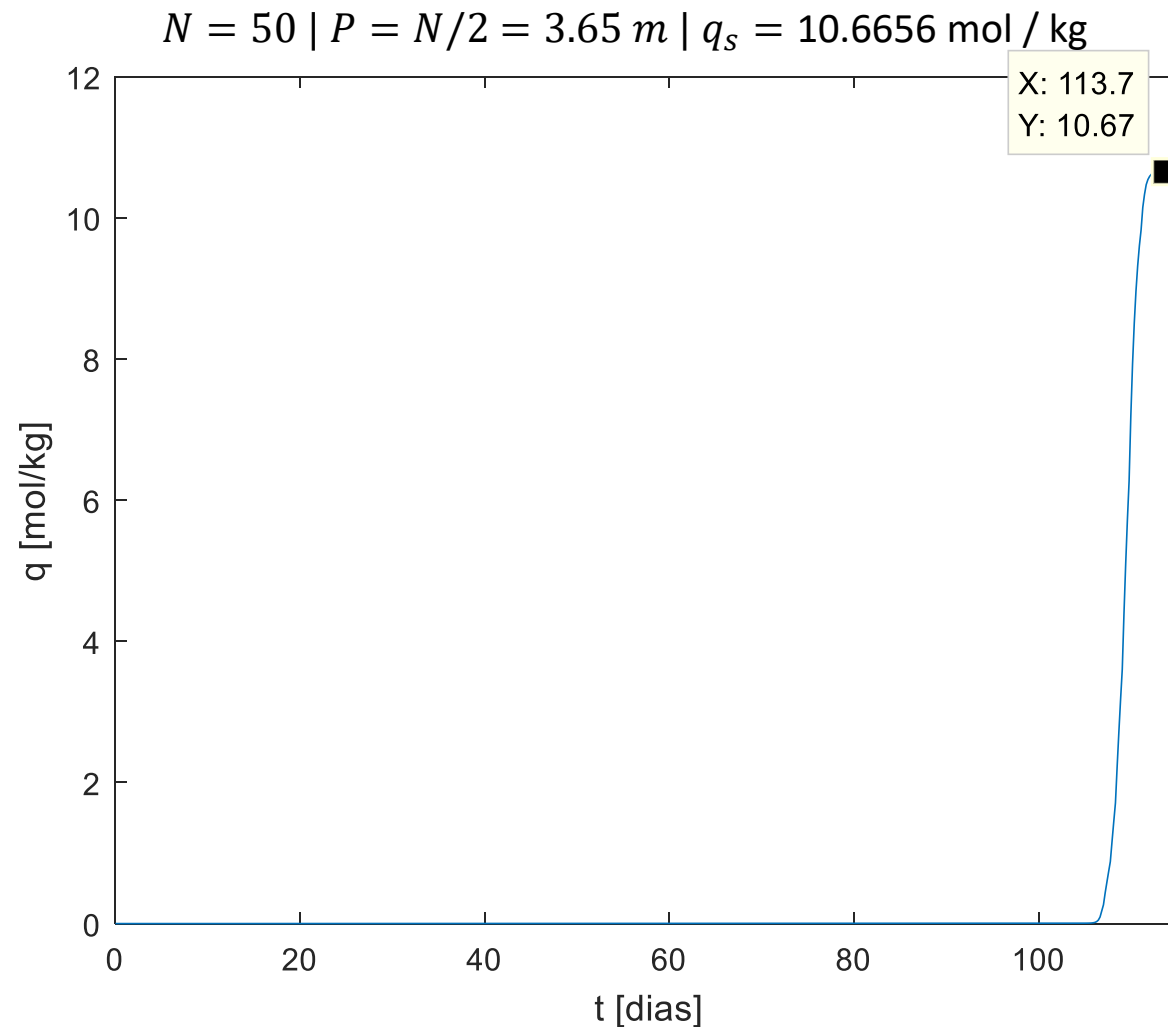
$0 < t < 20$

$C_0 = 18.8 \text{ mol/m}^3$

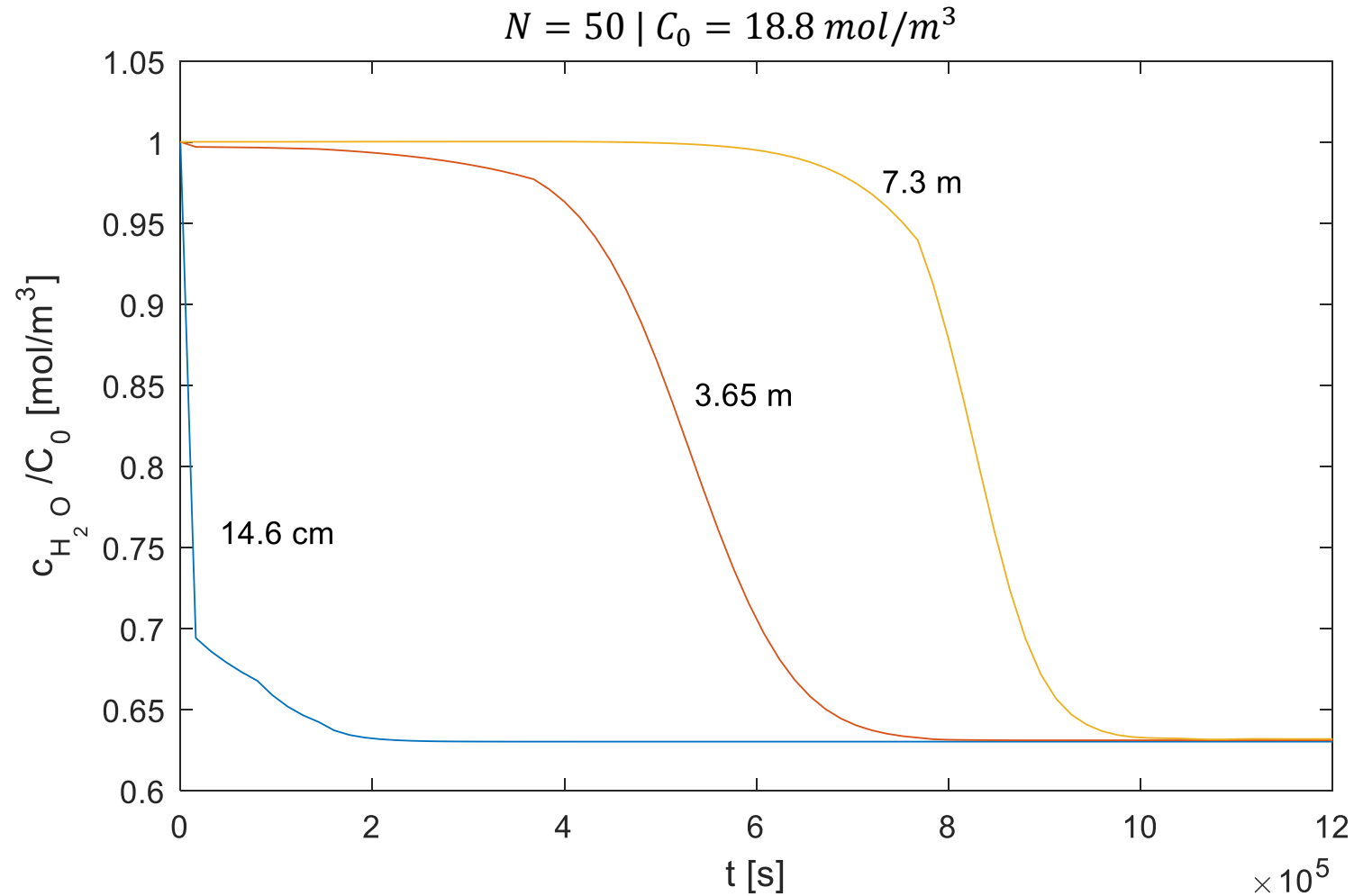
Concentração de Água no Gás



Água adsorvida por kg de material adsorvente

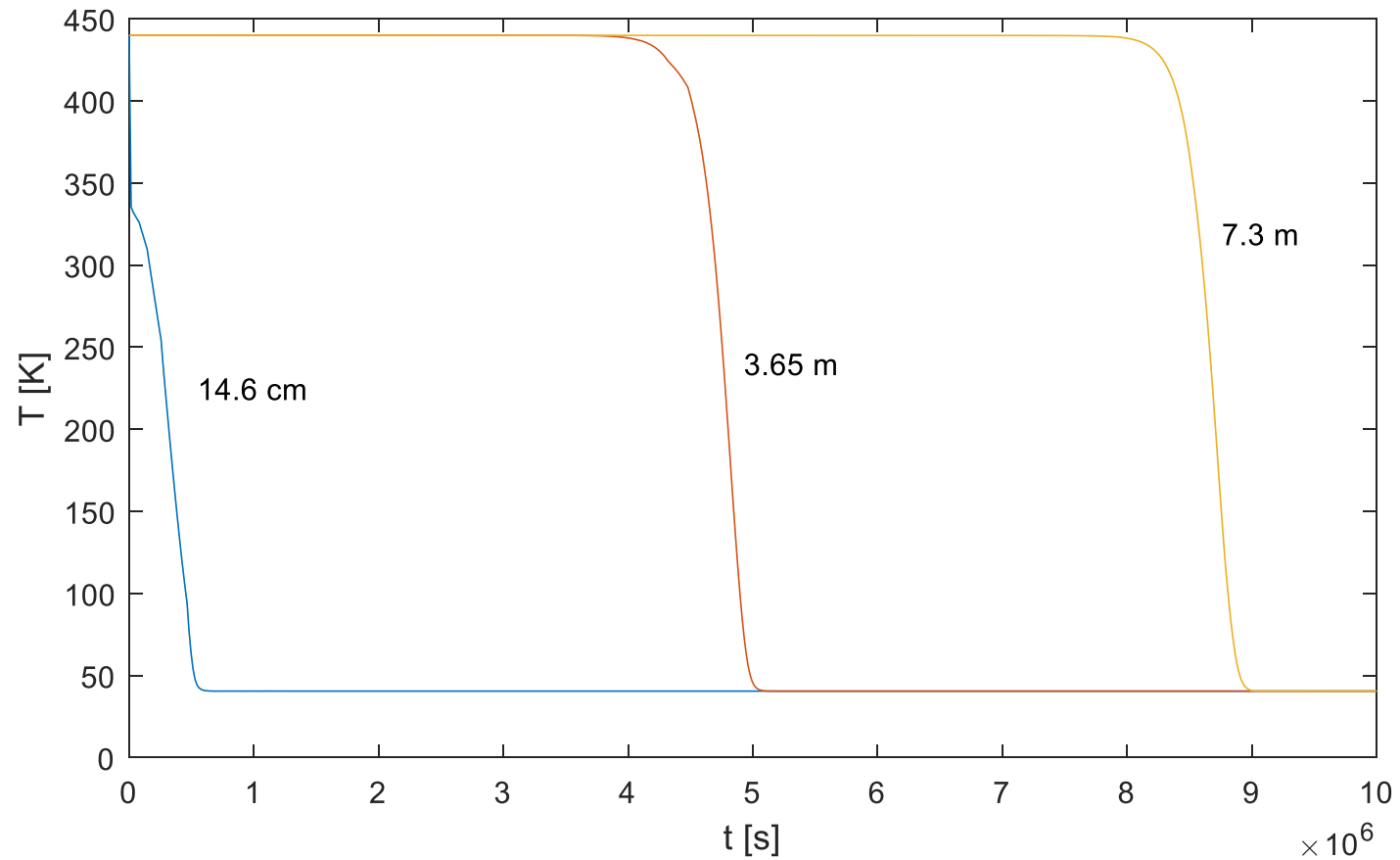


Curva de Ruptura



Variação de Temperatura

$N = 50 \mid T_0 = 440K$



Obrigada